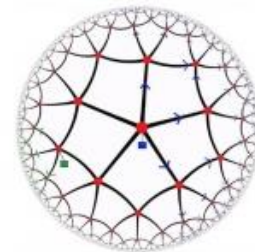


Tensor Networks, Entanglement, and Geometry

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Harvard and Brandeis

1607.05753 with John McGreevy



It from Qubit

Simons Collaboration on
Quantum Fields, Gravity and Information

Is quantum many-body physics “hard”?

- Experimental answer: obviously yes!
- Complexity answer: yes, because we can encode hard problems into quantum many-body physics – BQP, QMA, NP, glasses, and all that
- Hopeful theorist’s answer: certainly in some cases, but perhaps not in many (still poorly understood) cases of physical interest?
- By “easy” I mean: **a problem that can be solved in polynomial time on a classical computer**, *caveats: physical intuition/understanding?, unfavorable polynomial scaling?*

Two potentially “easy” classes?

- Ground states or thermal states of quantum field theories
 - Why? Renormalization group structure of entanglement → efficient tensor network representations (for regulated field theory)
[Vidal “MERA”, BGS-McGreevy “s-sourcery”, ...]
- Non-equilibrium steady states
 - Why? Coupling to environment → decoherence, low entanglement, local thermal equilibrium [Prosen-Znidaric, many others ...]
 - **Today:** a systematic quantum information approach to construct efficient classical representations of such states (and a connection to gravity)
[BGS-McGreevy 1607.05753,
Mahajan-Freeman-Mumford-Tubman-BGS 1608.05074]

Thermal states

$$\rho = \frac{e^{-H/T}}{Z}$$

$$H = \sum_x H_x$$

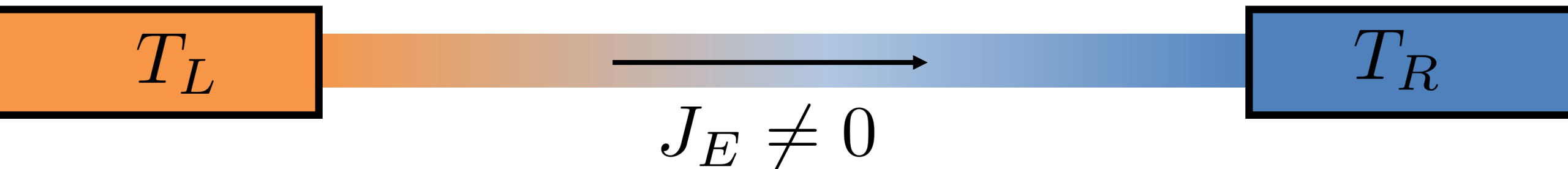
Local in space, few-body, translation invariant

- Typically have short-range correlations (except classical critical points)
- Area law for mutual information [Wolf-Cirac-Hastings-Verstraete]
- Always “trivial” for sufficiently high temperature

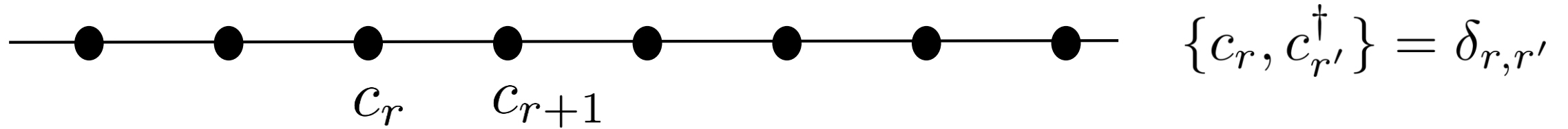
“Hydrodynamic” states

$$\rho \propto \exp \left(- \sum_x \frac{H_x}{T_x} + \dots \right)$$

- Local thermal equilibrium $\xi \nabla T_x \ll T_x$
- Generically carry currents, e.g. charge and heat



Non-interacting fermions



$$H = -w \sum_r \left[c_r^\dagger c_{r+1} + c_{r+1}^\dagger c_r \right]$$

solve using momentum eigenstates

$$H = \sum_k \epsilon_k c_k^\dagger c_k \quad \epsilon_k = -2w \cos k$$

thermal physics

$$\langle c_r^\dagger c_{r'} \rangle_T \sim e^{-|r-r'|/\xi} \quad \xi = a/T$$

Fermion thermal circuit

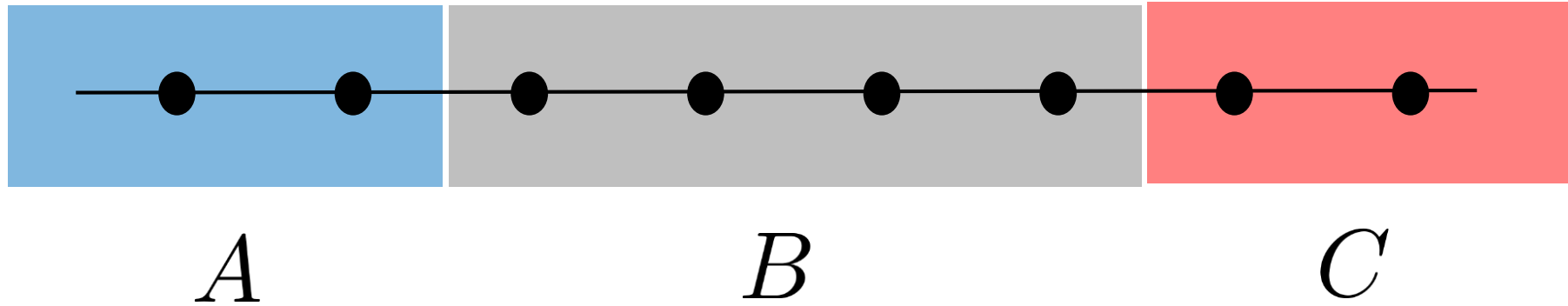
- Intuitively, the fermion thermal state should be easily preparable
- Construction: [\[BGS-McGreevy 1607.05753\]](#)
 - Introduce second copy of system
 - Find two-copy gapped Hamiltonian whose ground state is a purification of the original thermal state
 - Adiabatically evolve down from infinite temperature

$$H_T = \sum_k \left\{ (1 - 2f_k)c_k^\dagger c_k - (1 - 2f_k)\tilde{c}_k^\dagger \tilde{c}_k - 2\sqrt{f_k(1 - f_k)}(c_k^\dagger \tilde{c}_k + \tilde{c}_k^\dagger c_k) \right\}$$

interaction range in real space is the thermal length

$$H(\eta) = H_{T/\eta} \quad \text{family of Hamiltonians, gapped, bounded range, infinite T ground state is product}$$

What about the fermion system generalizes?



$$S(A) = -\text{Tr}(\rho_A \log \rho_A)$$

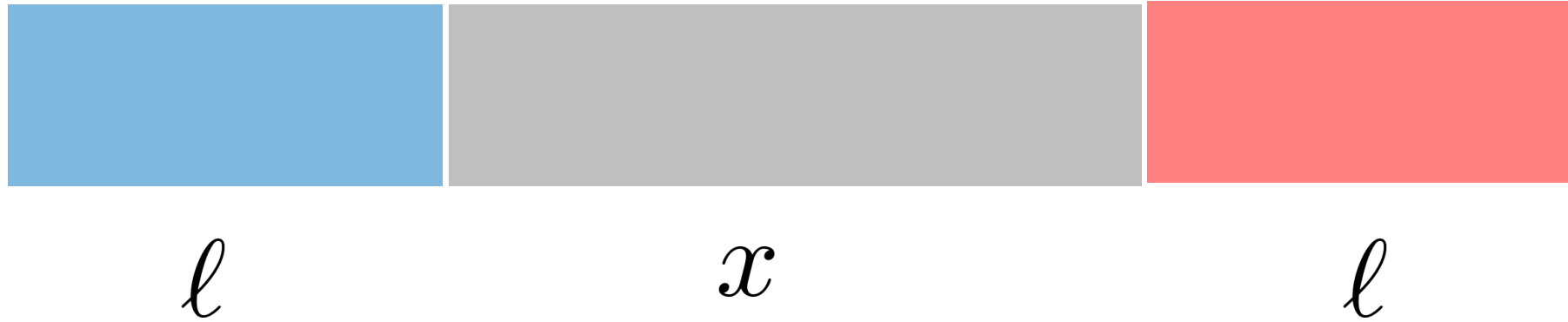
$$I(A : C|B) = S(AB) + S(BC) - S(B) - S(ABC)$$

conditional mutual information

$$I(A : C|B) \approx 0 \quad \text{if B is large enough}$$

[BGS-McGreevy 1607.05753]

Small CMI generalizes: conformal field theory



$$I(A : C|B) = \frac{c}{3} \log \left(\frac{\sinh^2[\pi T(\ell + x)]}{\sinh[\pi T x] \sinh[\pi T(2\ell + x)]} \right) \quad \text{exact for any CFT in 1+1d}$$

$$I(A : C|B) \sim ce^{-\pi T x} \quad \text{large } x$$

[BGS-McGreevy 1607.05753]

Reconstruction from small CMI

- When the conditional mutual information is zero, there is an exact reconstruction map that depends only on ρ_{AB} and ρ_{BC}

$$\rho_{ABC} = \rho_{BC}^{\frac{1}{2}} \rho_B^{-\frac{1}{2}} \rho_{AC} \rho_B^{-\frac{1}{2}} \rho_{BC}^{\frac{1}{2}} \quad [\text{Petz}]$$

- When the conditional mutual information is small, there is an approximate reconstruction map that depends only on ρ_{AB} and ρ_{BC}
[Fawzi-Renner, Sutter-Fawzi-Renner, Wilde...]
- Essentially, we are gluing local together to produce the global state

very much like hydrodynamics

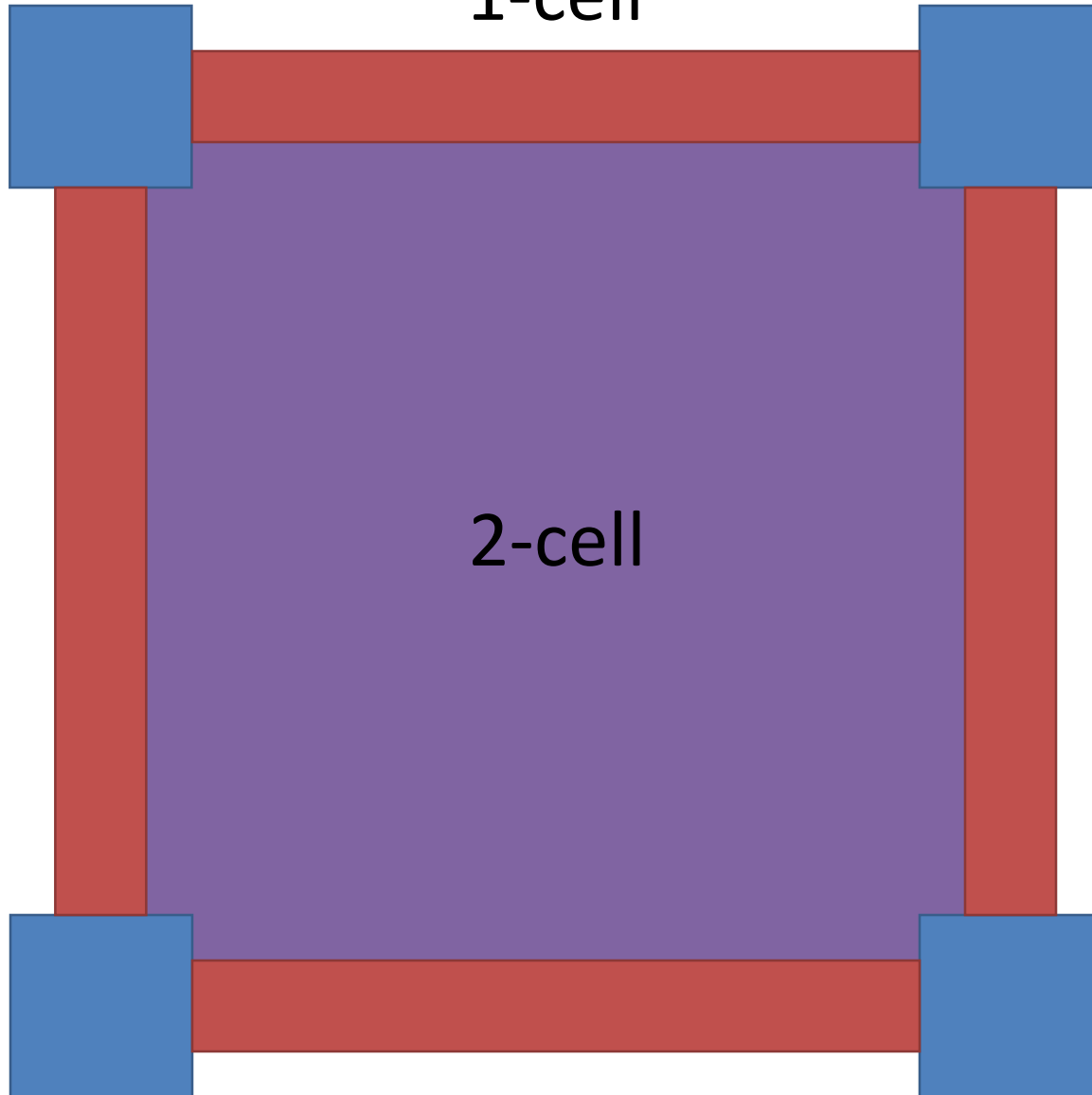
General entropic argument

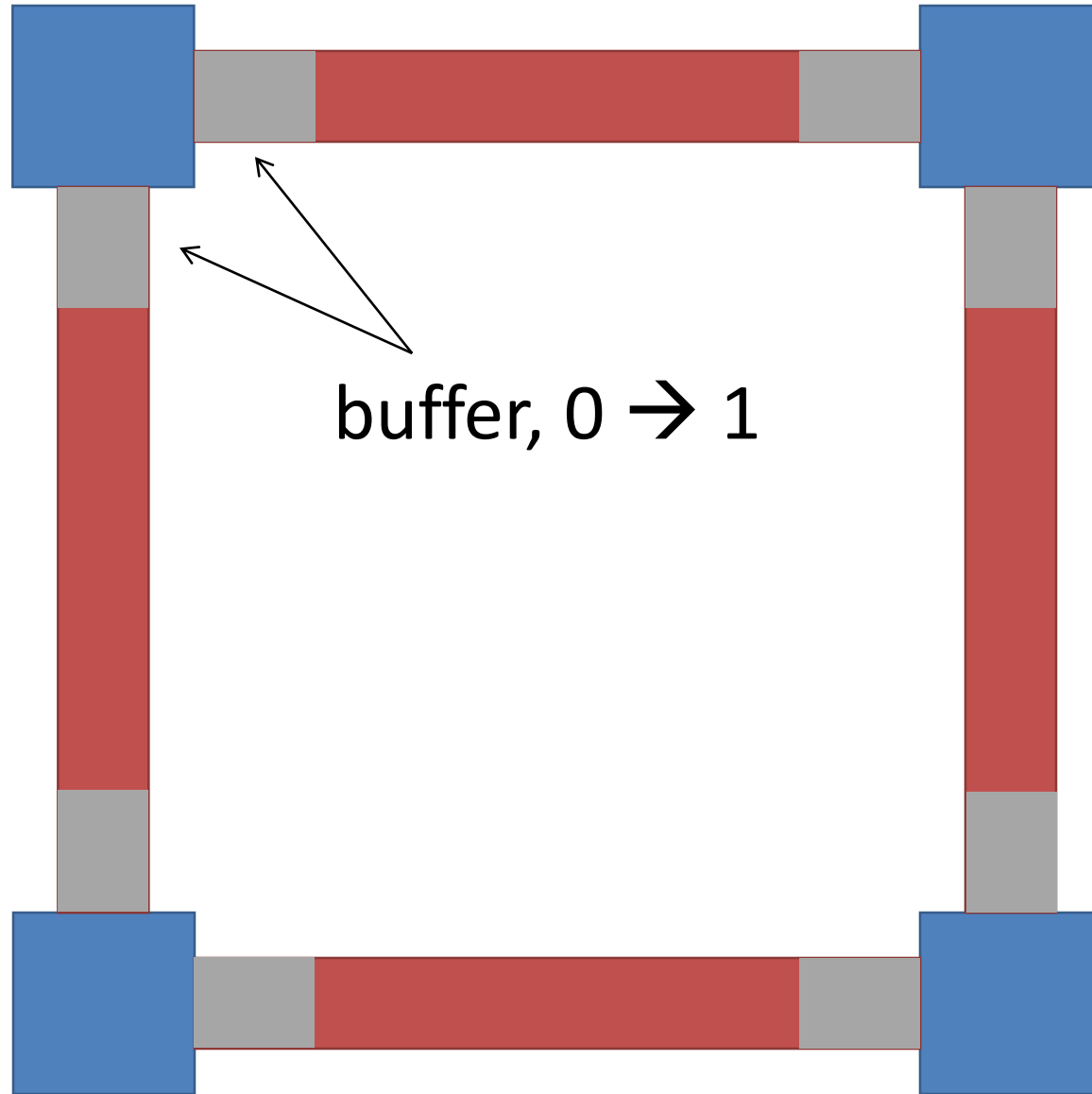
- For many quantum field theories (and quantum lattice models) at non-zero temperature, we expect the CMI of appropriate regions to approximately vanish
- Entropy $S = [\text{Volume Terms}] + [\text{Area Terms}] + [\text{Small corrections}]$
- Can be explicitly checked for a wide variety of models and is always valid at sufficiently high temperature, can fail for topological models e.g. 4d toric code / 2-form gauge theory
- **Main result: Combining CMI reconstruction results with a cellular construction shows that suitable thermal states can be efficiently prepared by short-range quantum channel** [BGS-McGreevy 1607.05753]

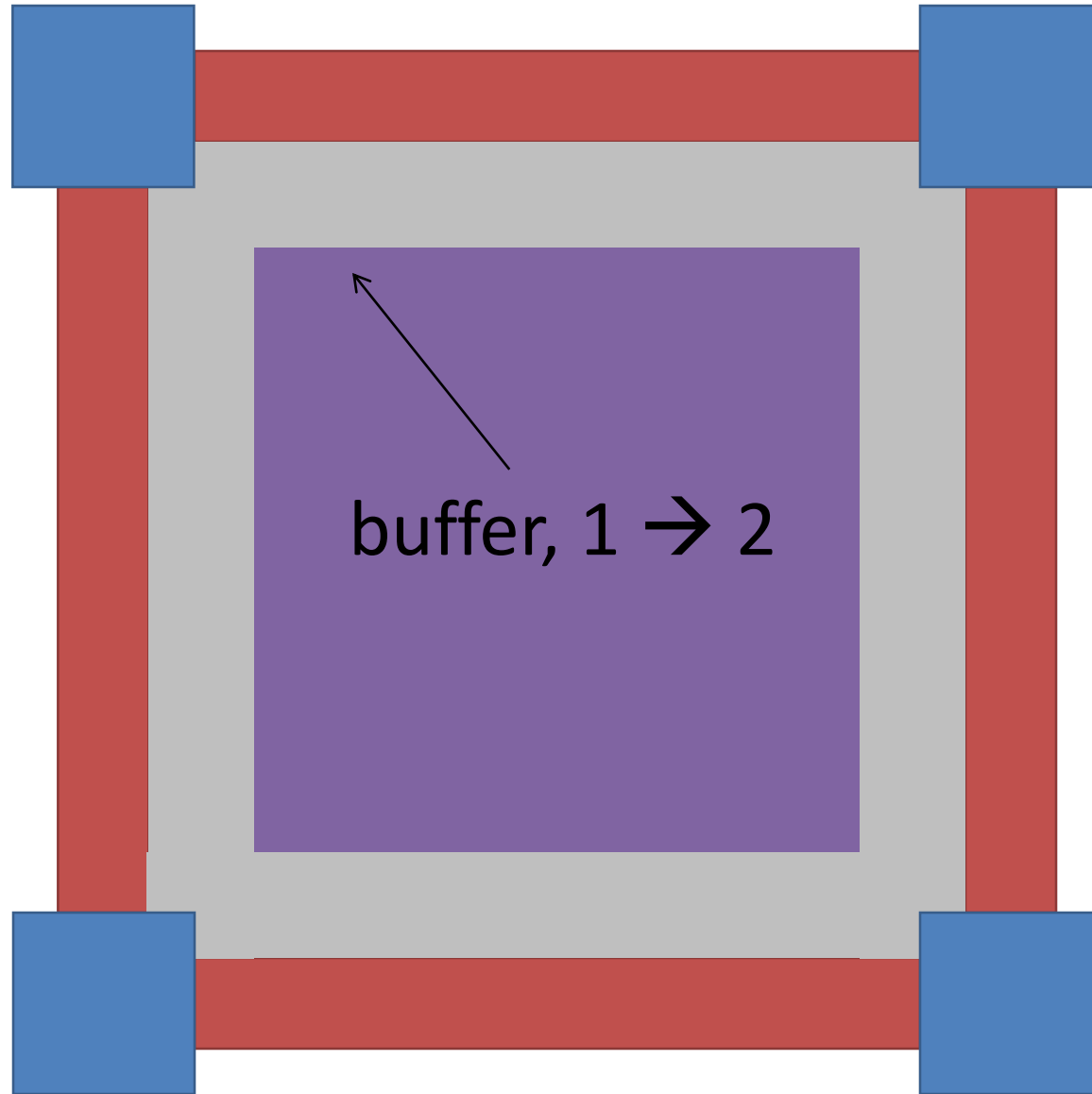
0-cell

1-cell

2-cell

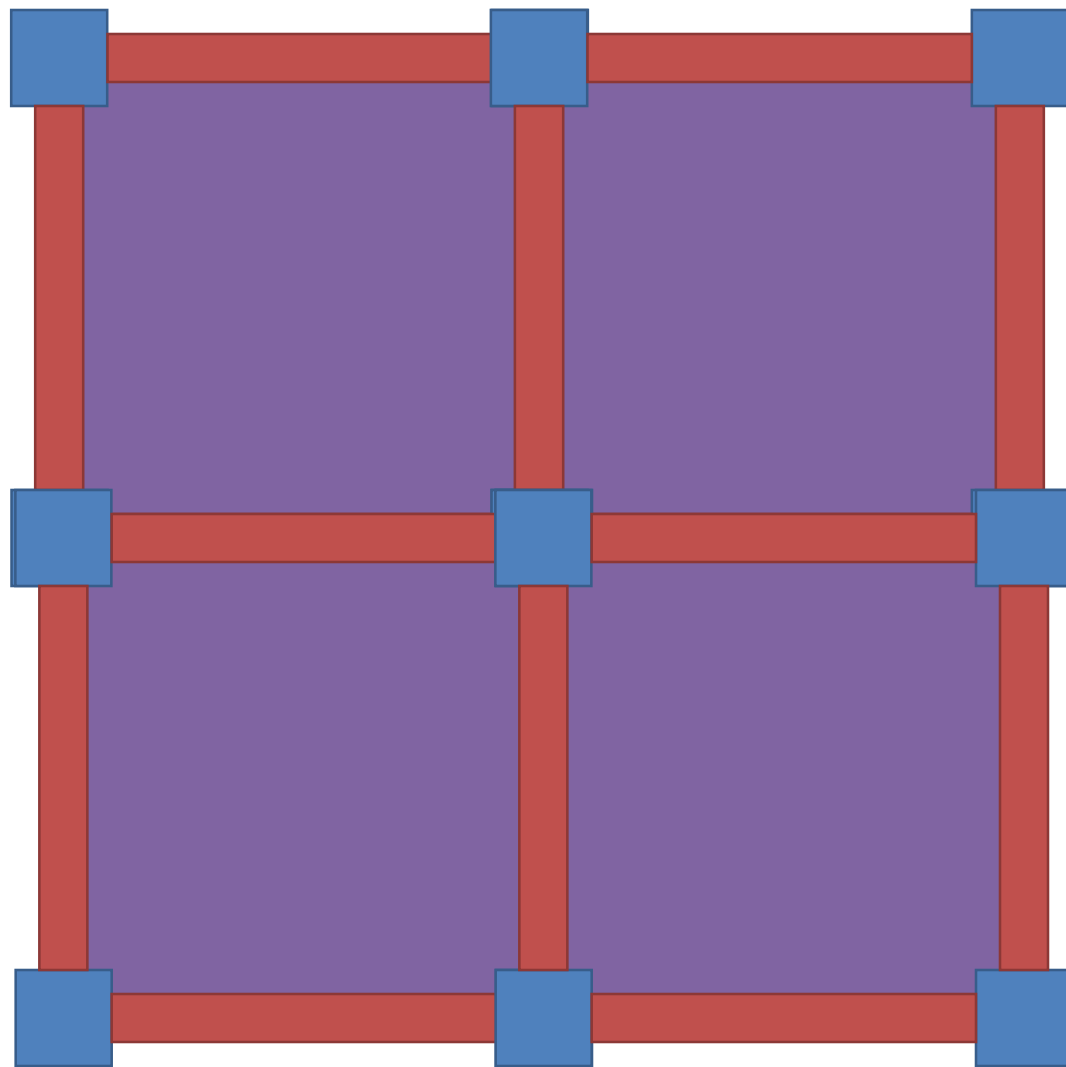




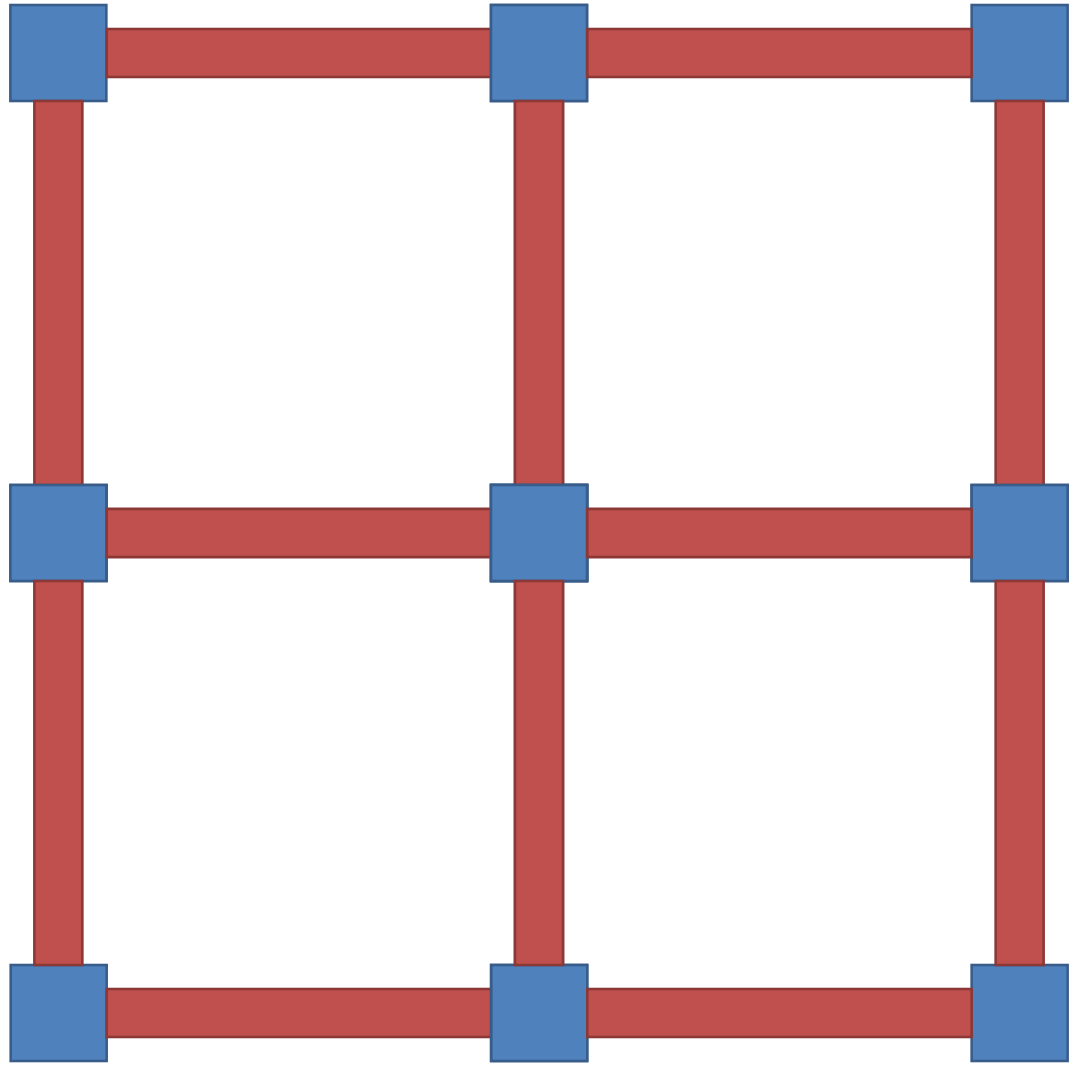


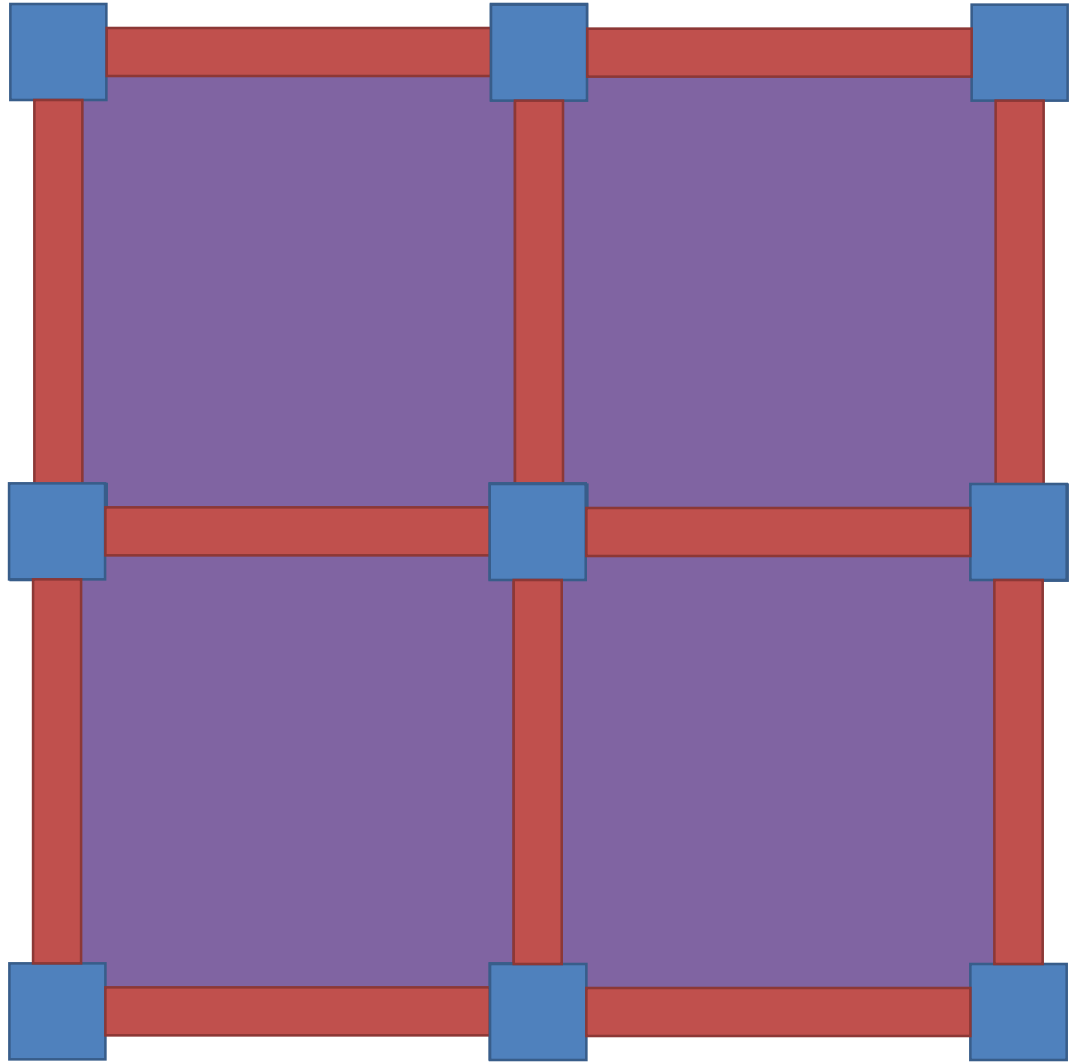
buffer, 1 → 2

Tile the plane with these unit cells:









Bubble of nothing array



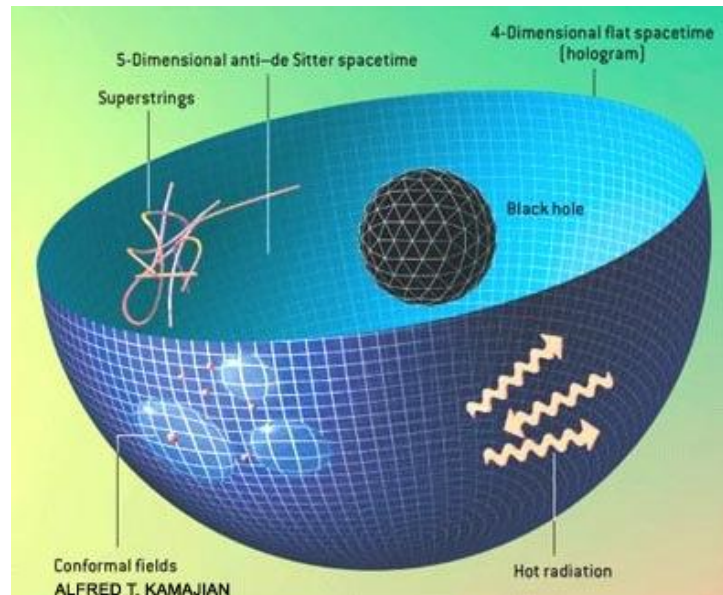
Main result

- Given: a state of n qubits with vanishing CMI for appropriate regions
- Result: there exists a pure state on a doubled system which is (approximately) the unique gapped ground state of a local Hamiltonian and which is adiabatically connected to a product state

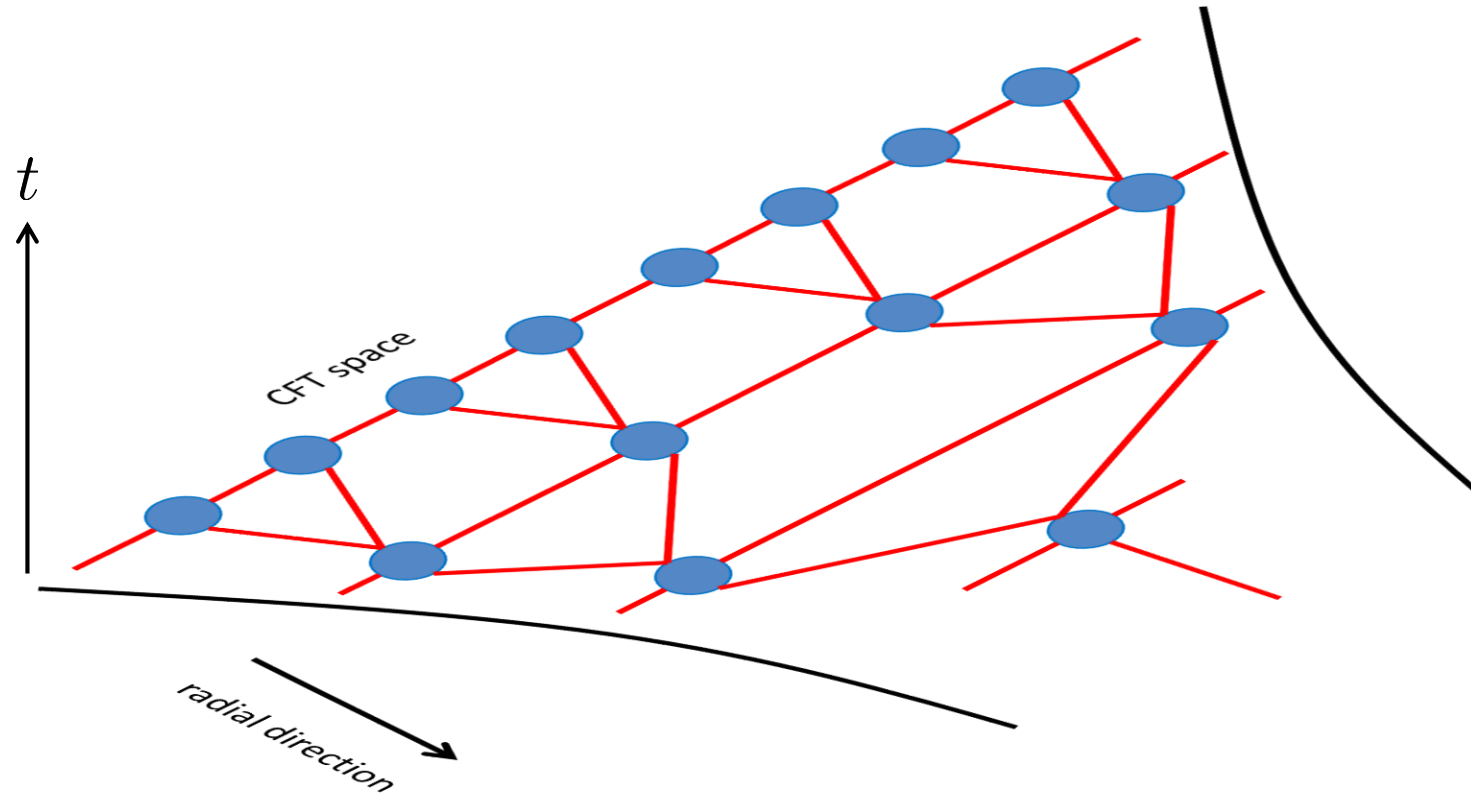
[BGS-McGreevy 1607.05753]

Connections to geometry and gravity

- AdS/CFT or holographic duality asserts: certain quantum field theories *without gravity* are exactly equivalent to certain quantum gravities in a higher dimension [Maldacena]
- Emergent direction is associated with energy scale in the field theory, an “RG direction”

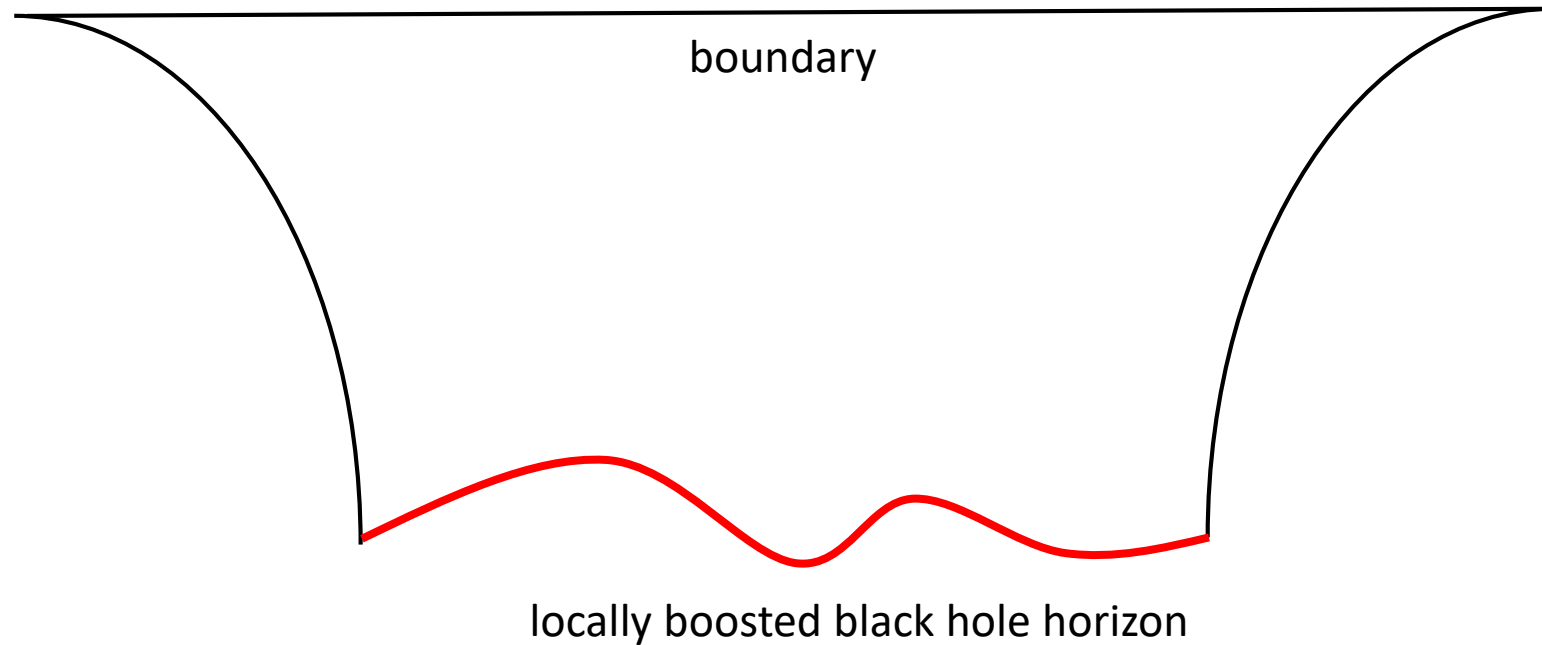


Spacetime as a tensor network [BGS]



Scale invariant CFT ground state $\rightarrow ds_{\text{AdS}}^2 = \frac{\ell^2}{r^2} (dr^2 - dt^2 + d\vec{x}^2)$

Hydrodynamic states in AdS/CFT



[Bhattacharyya-Hubeny-Minwalla-Rangamani]

Matching the tensor network and black hole descriptions

- AdS/CFT: thermal states of CFTs with holographic duals have approximately vanishing CMI [[BGS-McGreevy 1607.05753](#)]
- Remains true for hydrodynamic CFT states dual to hydrodynamic black holes [[BGS-Hubeny coming soon](#)]
- In both cases, we have a patch description of the state
 - Thermal scale chunks of the density matrix – local currents
 - Thermal scale patches of the black hole horizon – local boosts

Summary and work-in-progress

- Main result: efficient quantum information based representation of thermal states for a wide variety of quantum phases of matter
- Does the representation have an RG structure on scales smaller than the thermal scale?
- Can we derive hydrodynamic-like equations for the representation?
- Can we use the representation to make new calculations in strongly interacting systems?

THANK YOU